

Orientation of single crystals using linear approximations to NMR transits

Young-Sik Kye, Xingang Zhao, Gerard S. Harbison*

Department of Chemistry, University of Nebraska at Lincoln, Lincoln, NE 68588-0304, USA

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Abstract

We have derived analytical expressions for determining the orientation of high-symmetry single crystals from line-crossings in a single rotation plot. We demonstrate the utility of the method using the strontium-87 resonance in strontium nitrate. Employing our new method, which we call orientation of single crystals using linear approximations to NMR transits (OSCULANT), in combination with fourth-order perturbation theory, we obtain a highly accurate value for the quadrupole coupling constant, and an estimate for the chemical shielding anisotropy.

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1. Introduction

The most general method for determining magnitudes and orientations of internal NMR interactions uses single crystals to give well-resolved signals, whose orientation dependence can be analyzed to determine the secular part of, or higher-order perturbation contributions to, the nuclear spin Hamiltonian. In the standard method [1,2] the crystal is rotated in small increments about three successive, mutually orthogonal axes. Each rotation gives (for a secular, second rank interaction) two diagonal and one off-diagonal element of the Hamiltonian; three rotations thus provide, with some redundancy, all of the elements of the (necessarily symmetric) secular tensor. While more intellectually exciting methods have been proposed [3] using combinations of sample-hopping and multidimensional NMR, single crystal methods remain the mainstay of full tensor determination.

The crystal orientation for each of these three rotation plots is ordinarily obtained by X-ray diffraction or optical goniometry. Both methods require additional effort and equipment. Moreover, fixing the relative orientation of X-ray and NMR goniometers, on the one hand, and the orthogonality of the three rotation axes, on the other, remains the principal source of systematic error in single crystal studies. While intelligent design and careful machining [4] can minimize these errors, they cannot eliminate them.

In most crystalline systems, however, there are typically more than one chemically equivalent, magnetically distinct molecular species in the unit cell. The symmetry elements that generate this multiplicity are often regarded as a nuisance, since there is no general way to assign a particular signal to a given molecule. However, it is possible to use the symmetry properties of the unit cell itself, reflected in the NMR rotation plot, to provide absolute orientational information about the crystal in the frame of reference of the goniometer. This is only possible if there exist a multiplicity of NMR-active symmetry operations in the crystal. The most common such operation is the 2_1 screw operation, which (since NMR

* Corresponding author. Fax: +1 402 472 9402.

E-mail address: gerry@setanta.unl.edu (G.S. Harbison).

is oblivious to translations) is equivalent in NMR to a C_2 rotation about a crystallographic axis. Most monoclinic cells have a single such axis, which is insufficient fully to orient the crystal. However, many cells of orthorhombic, tetragonal or cubic symmetry have three mutually orthogonal 2_1 operations. In these cases, it is possible to determine the absolute orientation of the crystallographic a , b , and c axes, albeit usually not to distinguish between them (see, however, [5]).

Exploitation of such symmetry has until now been done in a fairly haphazard fashion. In early work [6], deviations of rotation plots from predicted patterns were used to correct small errors in crystal orientation. More recently, we determined [7] the unit cell orientation in an orthorhombic $Pbca$ crystal by iterative fitting to computed rotation plots. However, in the latter case the approximate crystal orientation was determined by optical goniometry, to generate computed rotation plots close enough to experiment to begin an iterative least squares fit.

In this paper, we present a simple method, which we call orientation of single crystals using linear approximation to NMR transits (OSULANT) to determine the orientation of the crystallographic axis system in the goniometer frame within a single rotation plot. The method works for any NMR Hamiltonian of even rank, and is therefore independent of the nature of the interaction or the crystal (so long as it contains the requisite number of symmetry operations). We will show that it gives excellent results from little more than visual inspection of the rotation plot, and presents a compelling argument for using crystals of high symmetry to determine NMR benchmark information. We demonstrate the method using a single set of rotation plot data for the ^{87}Sr resonances of $\text{Sr}(\text{NO}_3)_2$, using fourth-order perturbation theory to calculate the line positions, and then extract highly accurate values for the quadrupole coupling constant and the first estimate of a ^{87}Sr chemical shift anisotropy.

2. Methods

A single crystal of $\text{Sr}(\text{NO}_3)_2$ of approximate dimensions $0.8 \text{ mm} \times 0.5 \text{ mm} \times 0.3 \text{ mm}$ (Aldrich Chemical, St. Louis, MO, USA) was grown by slow evaporation from aqueous solution at $26^\circ \pm 2^\circ$. While at this temperature it is possible that either the anhydrous form or the dihydrate might crystallize, in our hands all of the crystals observed were cubic and therefore the anhydrous polymorph. The crystal was mounted in a singly tuned goniometer probe at a magnetic field of 7.187 T with a coil diameter of approximately 1.2 cm. All spectra were referenced to the ^{87}Sr signal from a freshly opened sample of SrO (Aldrich), which fell at 13.260013 MHz; the measured 90° pulse on this sample was 7 μs ; however,

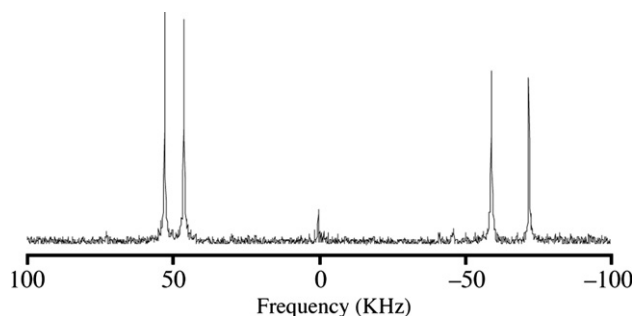


Fig. 1. ^{87}Sr NMR of a single crystal of strontium nitrate, centered at 13.246 MHz.

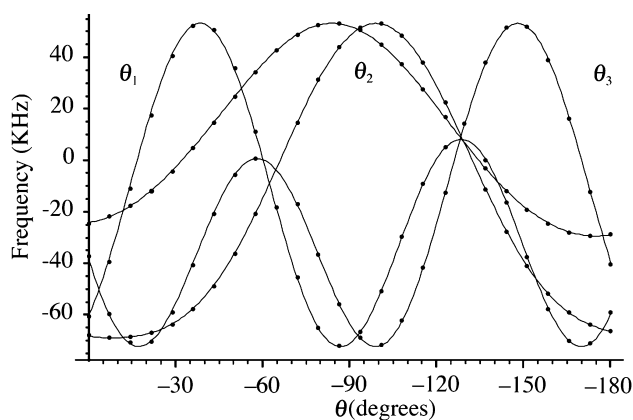


Fig. 2. Single crystal rotation plot for strontium nitrate, showing the three transits.

a small flip angle 1 μs pulse was used to excite the central transition of $\text{Sr}(\text{NO}_3)_2$. For each spectrum typically 15000 transients were signal-averaged with a relaxation delay of 0.5 s; a typical spectrum is shown in Fig. 1. The sample temperature during acquisition was $27 \pm 1^\circ\text{C}$. Data were collected in 7.2° increments for a full 180° rotation about the goniometer axis. The frequencies of the four identified strontium central transition resonances were obtained by a three-point fit to the peaks, and are plotted in Fig. 2. Note that the sign of the rotation in our NMR probe is negative, and therefore the data are plotted from 0° to -180° . The spectrum in Fig. 1 corresponds to $\theta = -86.8^\circ$ in the rotation plot.

2.1. Frames and conventions

Consider any interaction that can be described by a coupling tensor or combination of coupling tensors. We start with the principal axis frame (PAF) in which the tensor is diagonal. The transformation from PAF to crystallographic frame (XF), which in a crystal of orthorhombic or higher symmetry is defined by the crystallographic abc axes, is accomplished by the rotation matrix $R_C(\alpha_C, \beta_C, \gamma_C)$:

$$\begin{aligned}
R_C &= R_z(\alpha_C)R_y(\beta_C)R_z(\gamma_C) \\
&= \begin{pmatrix} \cos\alpha_C & -\sin\alpha_C & 0 \\ \sin\alpha_C & \cos\alpha_C & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \cos\beta_C & 0 & \sin\beta_C \\ 0 & 1 & 0 \\ -\sin\beta_C & 0 & \cos\beta_C \end{pmatrix} \\
&\quad \times \begin{pmatrix} \cos\gamma_C & -\sin\gamma_C & 0 \\ \sin\gamma_C & \cos\gamma_C & 0 \\ 0 & 0 & 1 \end{pmatrix}. \tag{1}
\end{aligned}$$

The four symmetry-related molecules in the unit cell are obtained from this single molecule by the rotation matrices $R_{S,0}$, $R_{S,x}$, $R_{S,y}$, and $R_{S,z}$:

$$\begin{aligned}
R_{S,0} &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}; \quad R_{S,x} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix}; \\
R_{S,y} &= \begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{pmatrix}; \quad R_{S,z} = \begin{pmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 1 \end{pmatrix}. \tag{2}
\end{aligned}$$

Transformation from this symmetry frame (SF) to the goniometer frame (GF) is accomplished by $R_G(\alpha_G, \beta_G, \gamma_G)$:

$$\begin{aligned}
R_G &= R_z(\alpha_G)R_y(\beta_G)R_z(\gamma_G) \\
&= \begin{pmatrix} \cos\alpha_G & -\sin\alpha_G & 0 \\ \sin\alpha_G & \cos\alpha_G & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \cos\beta_G & 0 & \sin\beta_G \\ 0 & 1 & 0 \\ -\sin\beta_G & 0 & \cos\beta_G \end{pmatrix} \\
&\quad \times \begin{pmatrix} \cos\gamma_G & -\sin\gamma_G & 0 \\ \sin\gamma_G & \cos\gamma_G & 0 \\ 0 & 0 & 1 \end{pmatrix} \tag{3}
\end{aligned}$$

and finally, rotation of the sample in the goniometer, by an angle θ , to the laboratory frame, to generate the actual rotation plot, about an axis which for historical reasons is defined as x , is accomplished by

$$R_L = R_x(\theta) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos\theta & -\sin\theta \\ 0 & \sin\theta & \cos\theta \end{pmatrix}. \tag{4}$$

The overall transformation from PAF to LF is the product:

$$R(\alpha, \beta, \gamma) = R_C(\alpha_C, \beta_C, \gamma_C)R_{S(0,x,y,z)}R_G(\alpha_G, \beta_G, \gamma_G)R_L(\theta). \tag{5}$$

And a tensor T in the lab frame is related to its diagonal form A_T by:

$$T = R^{-1}(\alpha, \beta, \gamma)A_T R(\alpha, \beta, \gamma). \tag{6}$$

2.2. Transits

Consider two molecules related by a π rotation about the y axis of the laboratory frame. If a second-rank sym-

metric coupling tensor relates the two molecules, the tensor T and T' will be identical except in the sign of the xz element, which will be opposite in sign in the two tensors. This element does not contribute to the NMR frequency to first order in perturbation theory, and to higher orders it enters only as even powers. This, the contribution to the NMR frequency of a second rank coupling tensor will be unaffected by a π_y rotation; and similar arguments can be made for a π_x rotation.

If the crystal containing these two molecules mounted in a random orientation in a goniometer, and is rotated orthogonal to \mathbf{z} , it is infinitely improbable that any of the crystallographic axes will by accident ever lie exactly parallel to \mathbf{z} ; however, each of these three axes must lie perpendicular to \mathbf{z} twice in any 2π rotation (Fig. 3). At the instant in the rotation where either a , b , or c axis lies in the xy plane, the NMR frequencies of the molecules related by a π rotation about the relevant crystallographic axis become twofold degenerate. In a rotation plot, this is visible as a point where all of the traces of individual NMR line cross pairwise. We label the a , b , or c axis transit points as θ_1 , θ_2 , and θ_3 .

In the present study, we rotated a crystal of $\text{Sr}(\text{NO}_3)_2$ about an axis perpendicular to the field, and labeled as x , and plotted the frequency of the central transition versus θ . The rotation plot is shown in Fig. 2. The three transits are indicated. Of considerable help in assigning these transits, and in distinguishing between symmetry-related transits and accidental degeneracies, is the knowledge that each curve transits each other exactly once per π rotation, and transits of both pairs of rotation plots must coincide exactly. Thus, accidental degeneracies in the plot in the region of -60° degrees can be discarded because the exact crossing points do not match. We find that accuracy of a few degrees in estimating the transit angle can be achieved by inspection, while linear intrapolation between the nearest frequencies generally gives an accuracy of better than a degree.

In the CF, the a axis obviously has coordinates $(1, 0, 0)$. In the lab frame, it has coordinates

$$a_L = R_L^{-1}(\theta)R_G^{-1}(\alpha_G, \beta_G, \gamma_G) \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}. \tag{7}$$

At $\theta = \theta_1$, the z component of a_L , $a_{L,z} = 0$ (since a now lies in the xy plane). This puts a constraint on the three Euler angles α_G , β_G , and γ_G . A second constraint is obtained at $\theta = \theta_2$ where $b_{L,z} = 0$, and a third at $\theta = \theta_3$ where $c_{L,z} = 0$. Inverting this set of three equations in α_G , β_G , and γ_G in terms of θ_1 , θ_2 , and θ_3 involves a considerable amount of algebra, and so we will give only the results.

$$\cos^2\beta_G = \frac{1 + \cot\theta_2 \cot\theta_3}{(\cot\theta_1 - \cot\theta_2)(\cot\theta_1 - \cot\theta_3)}, \tag{8}$$

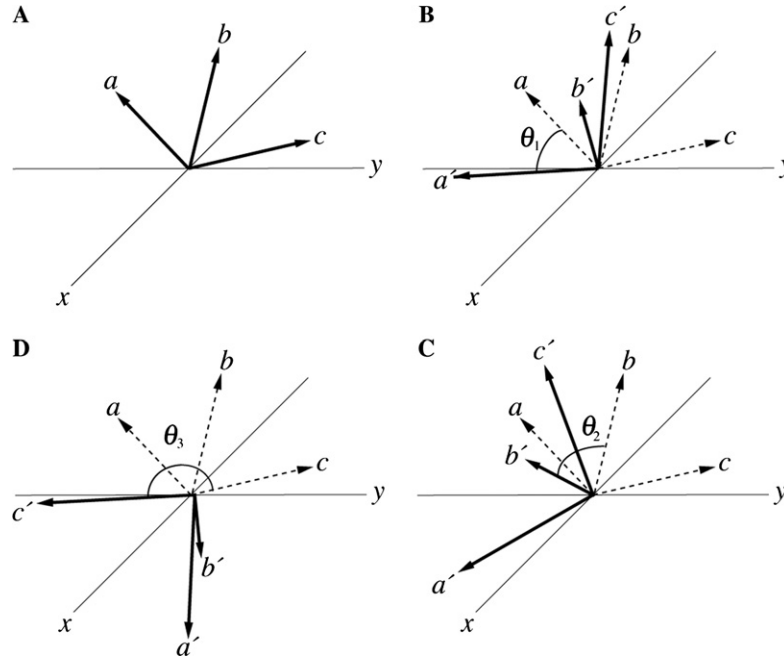


Fig. 3. A set of orthogonal crystal axes represented in the lab frame (A) with no rotation, and (B–D) rotated about x , by (B) θ_1 , so a' lies in the xy plane (C) θ_2 , so b' lies in the xy plane (D) θ_3 , so c' lies in the xy plane.

$$\sin^2 \gamma_G = \frac{\cot^2 \theta_1 (1 + \cot \theta_2 \cot \theta_3)}{\cot^2 \theta_1 + \cot \theta_1 (\cot \theta_2 + \cot \theta_3) - 1}, \quad (9)$$

$$\begin{aligned} \sin^2 2\alpha_G &= \frac{4(\cot \theta_1 - \cot \theta_2)(1 + \cot \theta_1 \cot \theta_2)(\cot \theta_1 - \cot \theta_3)(1 + \cot \theta_1 \cot \theta_3)}{(\cot \theta_2 - \cot \theta_3)^2 (1 - \cot^2 \theta_1 + \cot \theta_1 (\cot \theta_2 + \cot \theta_3))^2}. \end{aligned} \quad (10)$$

It is immediately apparent that these equations uniquely define β_G and γ_G only over the range 0 to $\pi/2$, and α_G only over the range $0 < \alpha_G < \pi/4$. However, the expressions for the three cotangents allow these ambiguities to be resolved:

$$\sin \gamma_G \tan \beta_G = -\cot \theta_1, \quad (11)$$

$$\cot \theta_2 = \frac{\sin \alpha_G \cos \beta_G \sin \gamma_G - \cos \alpha_G \cos \gamma_G}{\sin \alpha_G \sin \beta_G}, \quad (12)$$

$$\cot \theta_3 = \frac{\cos \alpha_G \cos \beta_G \sin \gamma_G + \sin \alpha_G \cos \gamma_G}{\cos \alpha_G \sin \beta_G}. \quad (13)$$

There are six possible ways to label the three transits as θ_1 , θ_2 , and θ_3 ; and six ways to label the crystallographic axes a , b , and c (albeit three of the possible labelings have opposite-handed coordinate systems).

2.3. Hamiltonian

By symmetry, the electric field gradient (EFG) tensor at the strontium site in $\text{Sr}(\text{NO}_3)_2$ must be axially symmetric, and directed along 111 or symmetry-equivalent directions. In the principal axis frame, whose unique axis coincides with the 111 axis, it is therefore given by

$$\mathbf{V} = \begin{pmatrix} -V_{zz}/2 & 0 & 0 \\ 0 & -V_{zz}/2 & 0 \\ 0 & 0 & V_{zz} \end{pmatrix}. \quad (14)$$

In the PAS frame, the Hamiltonian for the electric quadrupole coupling is:

$$\begin{aligned} \mathbf{H}_Q &= \left(\frac{eQ}{h(2I(2I-1))} \right) \mathbf{I} \cdot \mathbf{V} \cdot \mathbf{I} \\ &= \left(\frac{eQV_{zz}}{h(4I(2I-1))} \right) (3I_{z,PAS}^2 - \mathbf{I} \cdot \mathbf{I}), \end{aligned} \quad (15)$$

where $(eQV_{zz})/h$ is the so-called quadrupole coupling constant. The angular momentum vector may then be transformed into the laboratory frame:

$$I_{z,PAS} = I_z \cos \beta + I_x \sin \beta \cos \gamma + I_y \sin \beta \sin \gamma. \quad (16)$$

As is conventional, the laboratory frame carries no frame designation. The full laboratory Hamiltonian is then

$$\mathbf{H} = \mathbf{H}_Z + \mathbf{H}_Q + \mathbf{H}_\sigma = \gamma B_0 (1 - \sigma) I_z + \mathbf{H}_Q. \quad (17)$$

Assuming the chemical shift to be entirely secular, and treating the quadrupole interaction as a perturbation, we obtain the usual result that the odd-order terms in the perturbation expansion of the quadrupole Hamiltonian for the $1/2$ to $-1/2$ transition are zero. The second-order perturbation term becomes:

$$v_{1/2 \rightarrow -1/2}^{(2)} = \left(\frac{eQV_{zz}}{h} \right)^2 \frac{(1 - 10\cos^2 \beta + 9\cos^4 \beta)}{384\nu_0}, \quad (18)$$

where $\nu_0 = \gamma B_0$ is the nuclear Larmor frequency. This expression is identical to that of Volkoff [8] after substi-

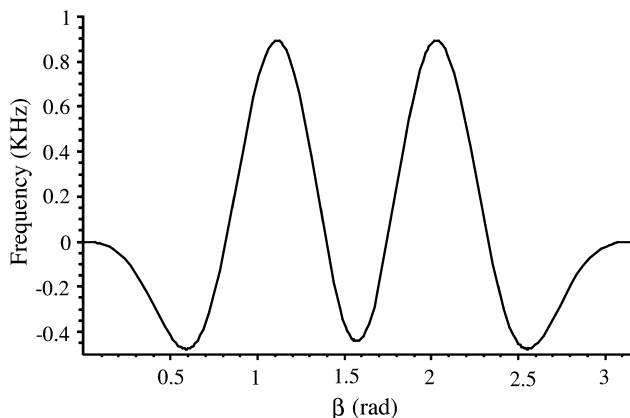


Fig. 4. Computed fourth-order quadrupole coupling for ^{87}Sr in strontium nitrate at 7 T, as a function of the angle β between EFG principal axis and the magnetic field.

tuting $\eta = 0$. Taking the expansion to fourth order by the method of Löwdin [9], we obtain

$$v_{-1/2-1/2}^{(4)} = -\left(\frac{eQV_{zz}}{h}\right)^4 \times \left(\frac{17 - 617\cos^2\beta + 2322\cos^4\beta - 2829\cos^6\beta + 1107\cos^8\beta}{884736v_0^3}\right). \quad (19)$$

This fourth-order term is plotted in Fig. 4, for a quadrupole coupling constant of 15.16 MHz and Zeeman frequency of 13.2 MHz. The $\cos 6\beta$ and $\cos 8\beta$ dependences are clear, and it can be seen that for these values of QCC and Zeeman frequency, the fourth-order contribution is of the order of 1 kHz, significantly larger than our final standard error of 560 Hz.

2.4. Data analysis

A trial rotation plot was calculated with an estimated QCC of 14.2 MHz, $\beta_C = \tan^{-1}\sqrt{2}$ and $\gamma_C = \pi/4$ as required by the space-group symmetry, and $(\alpha_G, \beta_G, \gamma_G) = (92.4^\circ, -82.2^\circ, -38.4^\circ)$, as determined using Eqs. (8)–(13) from $(\theta_1, \theta_2, \theta_3) = (-12.5^\circ, -93^\circ, -177^\circ)$. The trial plot was found to lie very close to the experimental data; using it, the four independent rotation plots were assigned, and the fit optimized by least-squares methods, varying the QCC, α_G , β_G , γ_G , and a correction to allow for misalignment of the x axis slightly away from perpendicular to the field. After optimization, the residual was found to contain both 2θ and higher-order dependences, and therefore, fourth-order corrections and a chemical shift anisotropy term were added (note that only the latter constitutes an additional degree of freedom; the former is merely an improvement in accuracy of the calculations). The fit was reoptimized, giving final

values of $(\alpha_G, \beta_G, \gamma_G) = (92.4^\circ, -82.6^\circ, -40.5^\circ)$, very close to the initial estimate from the transits. The errors given are 95% confidence limits, computed for each fit parameter from the total variance and a numerical estimate of the second partial derivative of the variance with respect to each parameter [10]. The goniometer axis misalignment was found to be 0.61° , while the optimized value of the QCC was 15.163 ± 0.018 MHz. The chemical shift with respect to SrO (sigma) was -428 ± 9 ppm, while the chemical shift anisotropy, $\sigma_{zz} - \sigma_i$, was -23 ± 28 ppm.

3. Discussion

The QCC obtained in this work is in fair agreement with the earlier, less precise measurement of 15.4 ± 0.2 MHz, obtained using spin-echo NMR of a strontium nitrate powder [11]. The slightly lower value obtained by us may be due to a difference in the measurement temperature, or it may be a result of the previous authors' use of η as a free parameter ($\eta = 0$ by symmetry in this system). The chemical shift anisotropy we obtain, while highly imprecise, is the first estimate we are aware of for ^{87}Sr (with the exception of systems such as SrO where it is zero by symmetry); it is certainly small and most likely negative, and should be compared with the value of -35 ppm measured for the more polarizable Pb^{2+} ion in the isomorphous $\text{Pb}(\text{NO}_3)_2$ crystal.

The present example was chosen because it is cubic, and therefore, the issue of assigning the individual a , b , and c axes does not arise; in addition, the orientation of the various NMR interactions in the crystallographic frame of reference is fixed by the local point symmetry. It might appear that in the more general case of an orthorhombic crystal with no point symmetry at the NMR nucleus, the inability to assign a , b , and c axes poses a problem. In fact, however, this is no more of a problem than if the axes have been assigned by some other means; because, even if the a , b , and c axes are known, the four symmetry-related molecules in the unit cell must still be assigned to four sets of NMR resonances. This can either be done using NMR interactions whose direction is known—dipolar splittings or dipolar linewidth—by using atoms which fall at special positions with respect to one of the crystal axes [5], or by comparison of tensor orientations with the local symmetry or with computed tensors. Once this assignment is done, of course, the three axes can also immediately be assigned. The general problem of assignment of chemically equivalent signals in high-symmetry crystals is usually a significant problem only where the chemical shielding anisotropies and local symmetry are low, as, for example, was the case for

sucrose [12]. In most cases, in practice, assignment is straightforward.

The OSCULANT method breaks down if the rotation plots contain insufficient information to orient the crystal, which can happen if a crystal axis either coincides with, or falls perpendicular to, the goniometer axis. In the former case, all NMR lines become twofold degenerate at all orientations, leaving insufficient information to determine all six tensor elements; in the latter, two transits occur at a single point in the rotation plot, leaving the orientation of the two axes within that plane undetermined. Mounting the crystal close to either of these limiting conditions causes the error in one or more of the fit parameters to become disproportionately large. In practice, we find that they can usually be avoided by making sure that the normals to all crystal faces are neither parallel nor perpendicular to the goniometer axis. Achieving a truly unsymmetrical mounting of the crystal is often surprisingly difficult, since the human brain seems to seek out symmetry even where it is being consciously avoided! Of course, an unfortunate choice is immediately apparent from the rotation plot, where the ideal is three well-spaced transits. And in practice, of course, statistics can be improved by conducting more than one differently mounted rotation, using a simultaneous refinement of tensor magnitudes and orientations while allowing the different crystal orientations to refine independently.

We hope this work convinces the interested reader that, if systematic errors due to crystal mounting can be eliminated by use of the symmetry-derived internal information in a rotation plot, remarkably precise tensors can be obtained, and estimates made of the chemical shielding of higher-spin nuclei even in the presence of very large quadrupolar interactions. This constitutes a strong argument for choosing crystals of orthorhombic or higher symmetry for single-crystal studies, since the disadvantage posed by the assignment problem is outweighed by the increased accuracy of the tensor determination.

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